

Notes on 2+1D Space Charge

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A useful formula is that if the full charge density is separable

$$\rho_{3D} = \rho(x, y)\lambda(z),$$

then the space charge equation $\nabla^2\phi = \rho_{3D}$ can be solved exactly by the infinite sum

$$\begin{aligned}\phi &= \sum_{n=0}^{\infty} (-1)^n \nabla^{-2(n+1)} \rho \lambda^{(2n)} \\ &= \sum_{n=0}^{\infty} (-1)^n \phi_{2D,n} \lambda^{(2n)},\end{aligned}$$

where we may calculate $\nabla^2\phi_{2D,0} = \rho$ and $\nabla^2\phi_{2D,n+1} = \phi_{2D,n}$ via Poisson solvers. The associated force law is

$$\mathbf{F} = \nabla\phi = \sum_{n=0}^{\infty} (-1)^n \begin{bmatrix} \nabla\phi_{2D,n} \lambda^{(2n)} \\ \phi_{2D,n} \lambda^{(2n+1)} \end{bmatrix}.$$

This series converges provided the repeated integrations transversely more than cancel the repeated derivatives longitudinally, which happens if ρ varies spatially faster than λ .

It is also true that a non-separable density may always be expanded into a sum of separable terms

$$\rho_{3D} = \sum_{n=0}^{\infty} \rho_n(x, y)\lambda_n(z),$$

each of which may be treated by the previous series and linearly combined. The convergence of this sum depends on the complexity of the charge distribution (the degree of longitudinal-transverse coupling), whereas the convergence of the other depends on the aspect ratio of detail in a long beam.

The traditional 2+1D formula corresponds in part to the $n = 0$ term of both series. Its force is proportional to

$$\mathbf{F} = \begin{bmatrix} \nabla\phi_{2D,0} \lambda \\ \phi_{2D,0} \lambda' \end{bmatrix}.$$

$\phi_{2D,0}$ is the solution to the usual transverse-only space charge problem. The transverse force term is exactly as expected, but the longitudinal term has λ' multiplied by the potential instead of a constant. Which other terms in the expansion make this resemble a constant times λ' is unclear, in fact it is not clear that a constant times λ' is a better approximation than this first term.